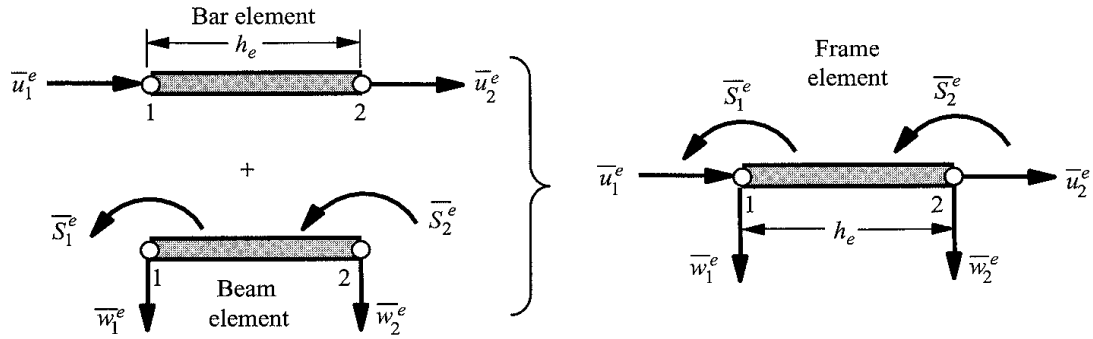


## Two-Dimensional Frame Elements



**Figure 5.4.1** Superposition of bar and beam element to obtain a frame element [degrees of freedom are referred to the element coordinate system  $(\bar{x}, \bar{y}, \bar{z})$ ].

A superposition of the bar element of Section 4.6 with the EBE of Section 5.2 or the Timoshenko beam element (RIE, CIE, or IIE) of Section 5.3 gives a frame element with three primary degrees of freedom ( $u$ ,  $w$ ,  $S$ ) per node (note that the transverse displacement  $v$  of Section 4.6 is now denoted by  $w$  to be consistent with Sections 5.2 and 5.3). When the axial stiffness  $EA$  and bending stiffness  $EI$  are elementwise constant, the superposition of the linear bar element with the IIE gives the following element equations (see Fig. 5.4.1):

$$[\bar{K}]^e \{\bar{\Delta}\}^e = \{\bar{F}\}^e \quad (5.4.1a)$$

or, in explicit form,

$$\frac{2EI}{\mu_0 h^3} \begin{bmatrix} \mu & 0 & 0 & -\mu & 0 & 0 \\ 0 & 6 & -3h & 0 & -6 & -3h \\ 0 & -3h & h^2(1.5 + 6\Lambda) & 0 & 3h & h^2(1.5 - 6\Lambda) \\ -\mu & 0 & 0 & \mu & 0 & 0 \\ 0 & -6 & 3h & 0 & 6 & 3h \\ 0 & -3h & h^2(1.5 - 6\Lambda) & 0 & 3h & h^2(1.5 + 6\Lambda) \end{bmatrix}^e \begin{Bmatrix} \bar{u}_1 \\ \bar{w}_1 \\ \bar{S}_1 \\ \bar{u}_2 \\ \bar{w}_2 \\ \bar{S}_2 \end{Bmatrix}^e = \begin{Bmatrix} \bar{F}_1 \\ \bar{F}_2 \\ \bar{F}_3 \\ \bar{F}_4 \\ \bar{F}_5 \\ \bar{F}_6 \end{Bmatrix}^e \quad (5.4.1b)$$

where

$$\begin{Bmatrix} \bar{F}_1 \\ \bar{F}_2 \\ \bar{F}_3 \\ \bar{F}_4 \\ \bar{F}_5 \\ \bar{F}_6 \end{Bmatrix}^e = \begin{Bmatrix} \bar{f}_1 \\ \bar{q}_1 \\ \bar{q}_2 \\ \bar{f}_2 \\ \bar{q}_3 \\ \bar{q}_4 \end{Bmatrix}^e + \begin{Bmatrix} \bar{Q}_1 \\ \bar{Q}_2 \\ \bar{Q}_3 \\ \bar{Q}_4 \\ \bar{Q}_5 \\ \bar{Q}_6 \end{Bmatrix}^e \quad (5.4.2a)$$

$$\bar{f}_i = \int_0^{h_e} f^e(\bar{x}) \psi_i^e(\bar{x}) d\bar{x} \quad (i = 1, 2), \quad \bar{q}_i = \int_0^{h_e} q^e(\bar{x}) \phi_i^e(\bar{x}) d\bar{x} \quad (i = 1, 2, 3, 4) \quad (5.4.2b)$$

and

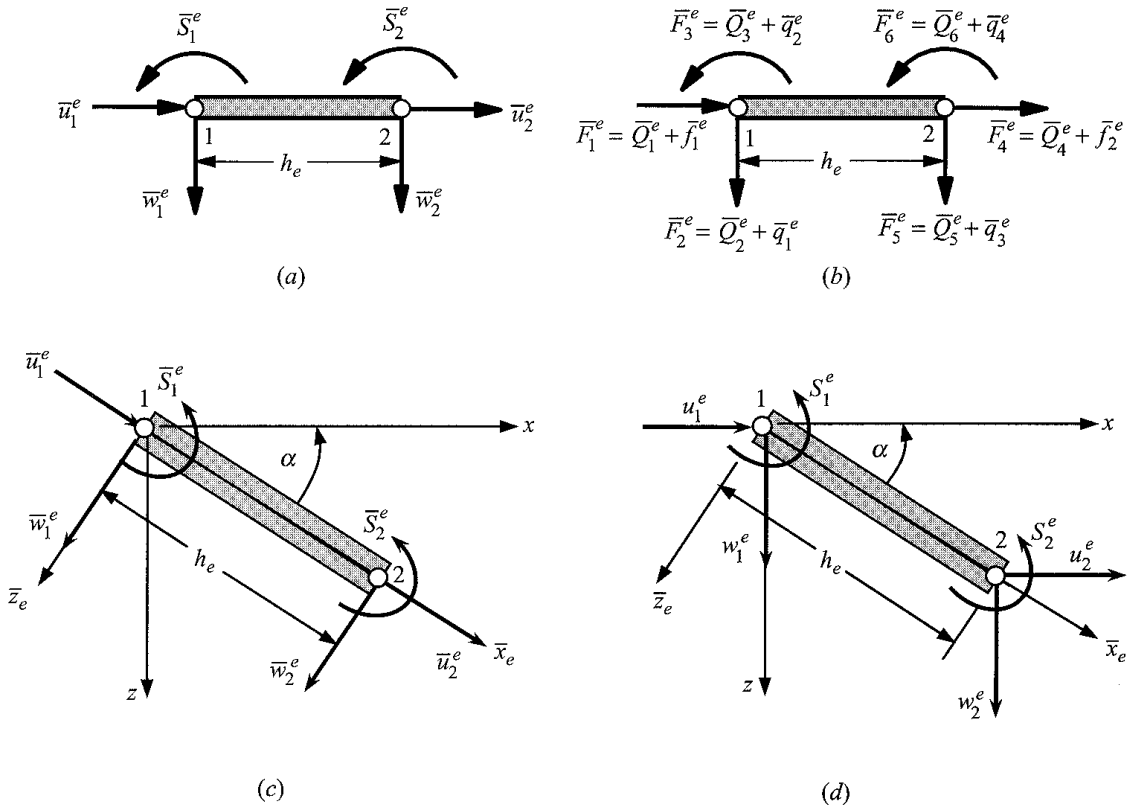
$$\mu = \frac{A\mu_0 h^2}{2I}, \quad \Lambda = \frac{EI}{GAK_s h^2}, \quad \mu_0 = 12\Lambda \quad (5.4.3)$$

In Eq. (5.4.2b),  $f^e$  denotes the distributed axial force,  $q^e$  the distributed transverse force,  $\psi_i^e$  the linear interpolation functions, and  $\phi_i^e$  the Hermite cubic interpolation functions. In the following paragraphs, we develop transformation relations to express the element equations (5.4.1b)—valid in the element coordinate system  $(\bar{x}, \bar{y}, \bar{z})$ —to the global coordinate system  $(x, y, z)$ .

The local coordinates  $(\bar{x}_e, \bar{y}_e, \bar{z}_e)$  of a typical element  $\Omega_e$  are related to the global coordinates  $(x, y, z)$  by [cf. Eq. (4.6.2)]

$$\begin{Bmatrix} \bar{x} \\ \bar{y} \\ \bar{z} \end{Bmatrix}^e = \begin{bmatrix} \cos \alpha & 0 & \sin \alpha \\ 0 & 1 & 0 \\ -\sin \alpha & 0 & \cos \alpha \end{bmatrix}^e \begin{Bmatrix} x \\ y \\ z \end{Bmatrix} \quad (5.4.4)$$

where the angle  $\alpha_e$  is measured clockwise from the global  $x$ -axis to the element  $\bar{x}_e$ -axis. Note that the  $y$  and  $\bar{y}_e$  coordinates are parallel to each other, and they are *out of* the plane of the paper (see Fig. 5.4.2). The same transformation relations hold for displacements  $(u, w)$  along the global coordinates  $(x, z)$  and displacements  $(\bar{u}, \bar{w})$  in the local coordinates



**Figure 5.4.2** (a) Generalized displacements. (b) Generalized forces. (c) Generalized displacements in the element coordinates. (d) Generalized displacements in the global coordinates.

$(\bar{x}, \bar{z})$ . Note that there is no displacement in the direction of the coordinate  $y$  (i.e.,  $v = 0$ ). However, there is a rotation about the  $y$ -axis, and it remains the same in both coordinate systems because  $y = \bar{y}$ . Note that rotation  $\theta$  is equal to  $-dw/dx$  in Euler–Bernoulli beam theory and it is equal to  $\Psi$  in Timoshenko beam theory. Hence, the relationship between  $(u, w, \theta)$  and  $(\bar{u}, \bar{w}, \bar{\theta})$  can be written as

$$\begin{Bmatrix} \bar{u} \\ \bar{w} \\ \bar{\theta} \end{Bmatrix}^e = \begin{bmatrix} \cos \alpha & \sin \alpha & 0 \\ -\sin \alpha & \cos \alpha & 0 \\ 0 & 0 & 1 \end{bmatrix}^e \begin{Bmatrix} u \\ w \\ \theta \end{Bmatrix}^e \quad (5.4.5)$$

Therefore, the three nodal degrees of freedom  $(\bar{u}_i^e, \bar{w}_i^e, \bar{S}_i^e)$  at the  $i$ th node ( $i = 1, 2$ ) in the  $(\bar{x}, \bar{y}, \bar{z})$  system are related to the three degrees of freedom  $(u_i^e, w_i^e, S_i^e)$  in the  $(x, y, z)$  system by

$$\begin{Bmatrix} \bar{u}_1 \\ \bar{w}_1 \\ \bar{S}_1 \\ \bar{u}_2 \\ \bar{w}_2 \\ \bar{S}_2 \end{Bmatrix}^e = \begin{bmatrix} \cos \alpha & \sin \alpha & 0 & & & \\ -\sin \alpha & \cos \alpha & 0 & & & \\ 0 & 0 & 1 & & & \\ & & & \cos \alpha & \sin \alpha & 0 \\ & & & -\sin \alpha & \cos \alpha & 0 \\ & & & 0 & 0 & 1 \end{bmatrix}^e \begin{Bmatrix} u_1 \\ w_1 \\ S_1 \\ u_2 \\ w_2 \\ S_2 \end{Bmatrix}^e \quad (5.4.6a)$$

or

$$\{\bar{\Delta}^e\} = [T^e]\{\Delta^e\} \quad (5.4.6b)$$

Analogously, the element force vectors in the local and global coordinate systems are related according to

$$\{\bar{F}^e\} = [T]^e\{F^e\} \quad (5.4.7)$$

Returning to Eq. (5.4.1a), we substitute the transformation equations (5.4.6b) and (5.4.7) into (5.4.1a) and obtain

$$[\bar{K}]^e [T]^e \{\Delta\}^e = [T]^e \{F\}^e$$

Premultiplying both sides with  $[T]^{-1} = [T]^T$ , we obtain

$$[T]^T [\bar{K}]^e [T]^e \{\Delta\}^e = \{F\}^e \quad \text{or} \quad [K]^e \{\Delta\}^e = \{F\}^e \quad (5.4.8)$$

where

$$[K]^e = [T]^T [\bar{K}]^e [T]^e, \quad \{F\}^e = [T]^T \{\bar{F}\}^e \quad (5.4.9)$$

Thus, if we know the element matrices  $[\bar{K}]^e$  and  $\{\bar{F}\}^e$  of an element  $\Omega_e$  in the local coordinate system  $(\bar{x}, \bar{y}, \bar{z})$ , the element matrices in the global coordinate system are obtained by (5.4.9).

Using  $[\bar{K}]^e$  and  $\{\bar{F}\}^e$  from Eq. (5.4.1b) in (5.4.9) and carrying out the indicated matrix multiplications, we arrive at the following element stiffness matrix  $[K]^e$  referred to the

global coordinates:

$$[K]^e = \frac{2EI}{\mu_0 h^3} \begin{bmatrix} \mu \cos^2 \alpha + 6 \sin^2 \alpha & (\mu - 6) \cos \alpha \sin \alpha & 3h \sin \alpha \\ (\mu - 6) \cos \alpha \sin \alpha & \mu \sin^2 \alpha + 6 \cos^2 \alpha & -3h \cos \alpha \\ 3h \sin \alpha & -3h \cos \alpha & h^2(1.5 + 6\Delta) \\ -(\mu \cos^2 \alpha + 6 \sin^2 \alpha) & -(\mu - 6) \sin \alpha \cos \alpha & -3h \sin \alpha \\ -(\mu - 6) \cos \alpha \sin \alpha & -(\mu \sin^2 \alpha + 6 \cos^2 \alpha) & 3h \cos \alpha \\ 3h \sin \alpha & -3h \cos \alpha & h^2(1.5 - 6\Delta) \\ -(\mu \cos^2 \alpha + 6 \sin^2 \alpha) & -(\mu - 6) \cos \alpha \sin \alpha & 3h \sin \alpha \\ -(\mu - 6) \sin \alpha \cos \alpha & -(\mu \sin^2 \alpha + 6 \cos^2 \alpha) & -3h \cos \alpha \\ -3h \sin \alpha & 3h \cos \alpha & h^2(1.5 - 6\Delta) \\ (\mu \cos^2 \alpha + 6 \sin^2 \alpha) & (\mu - 6) \cos \alpha \sin \alpha & -3h \sin \alpha \\ (\mu - 6) \cos \alpha \sin \alpha & \mu \sin^2 \alpha + 6 \cos^2 \alpha & 3h \cos \alpha \\ -3h \sin \alpha & 3h \cos \alpha & h^2(1.5 + 6\Delta) \end{bmatrix} \quad (5.4.10a)$$

$$\{F\}^e = \begin{Bmatrix} F_1 \\ F_2 \\ F_3 \\ F_4 \\ F_5 \\ F_6 \end{Bmatrix}^e = \begin{Bmatrix} \bar{F}_1 \cos \alpha - \bar{F}_2 \sin \alpha \\ \bar{F}_1 \sin \alpha + \bar{F}_2 \cos \alpha \\ \bar{F}_3 \\ \bar{F}_4 \cos \alpha - \bar{F}_5 \sin \alpha \\ \bar{F}_4 \sin \alpha + \bar{F}_5 \cos \alpha \\ \bar{F}_6 \end{Bmatrix}^e \quad (5.4.10b)$$

which is the element force vector referred to the global coordinates.